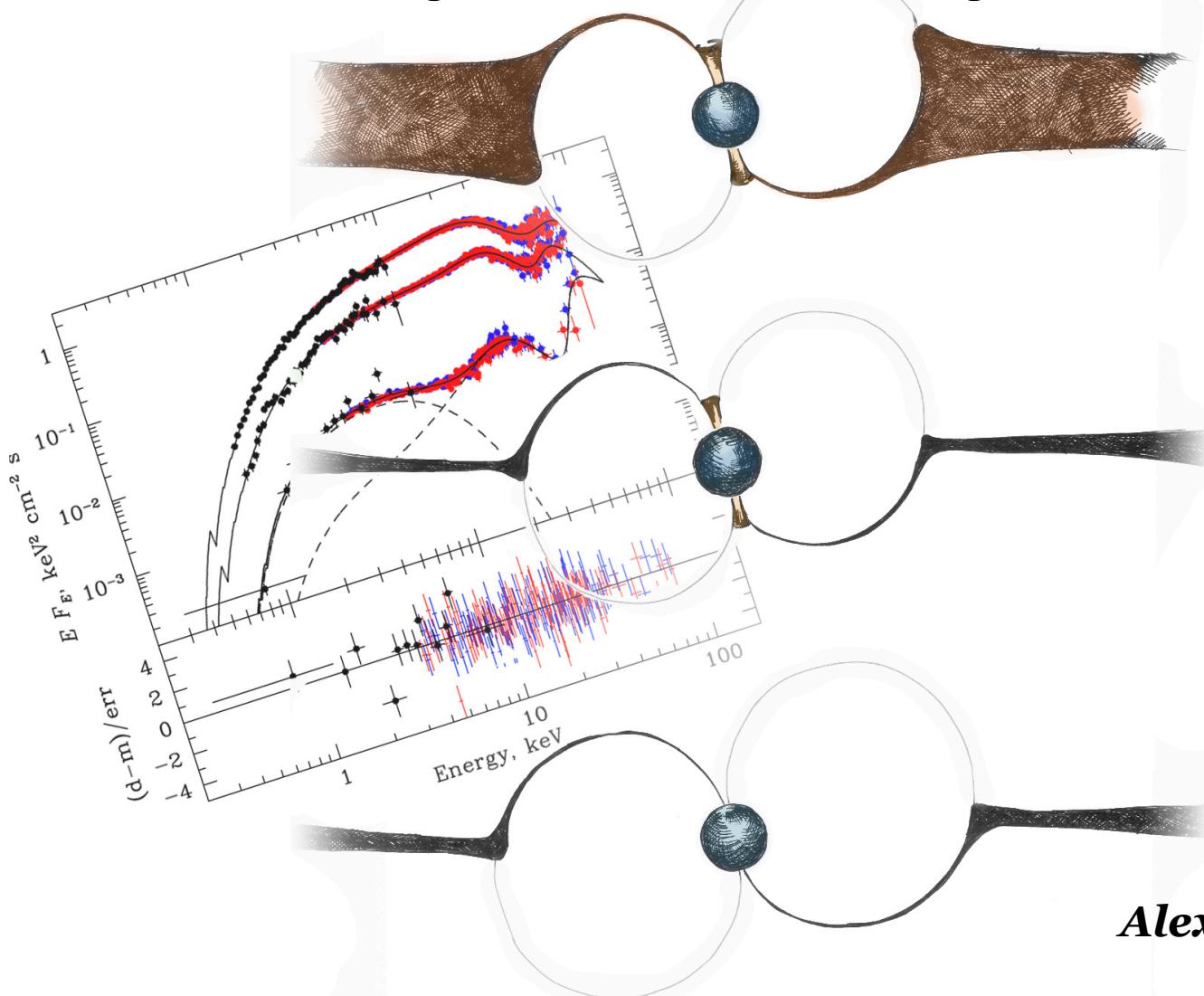


X-ray pulsars at very low luminosity state



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Victor Doroshenko
Alexander Lutovinov
Juri Poutanen*



Universiteit
Leiden



Главная (Пулковская)
астрономическая
обсерватория РАН

X-ray pulsar

Rotating Neutron Star in binary systems

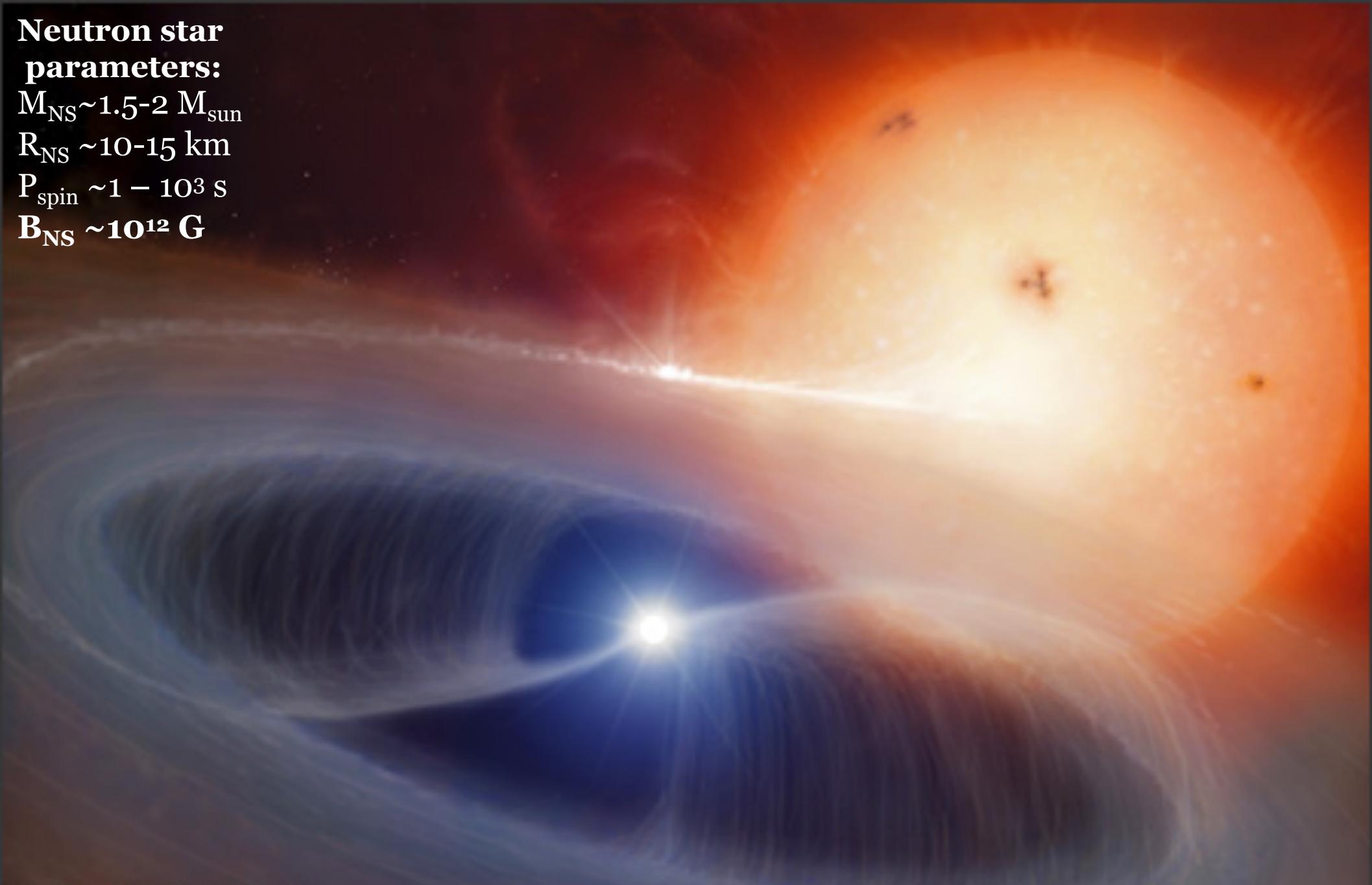
**Neutron star
parameters:**

$M_{NS} \sim 1.5 - 2 M_{\text{sun}}$

$R_{NS} \sim 10 - 15 \text{ km}$

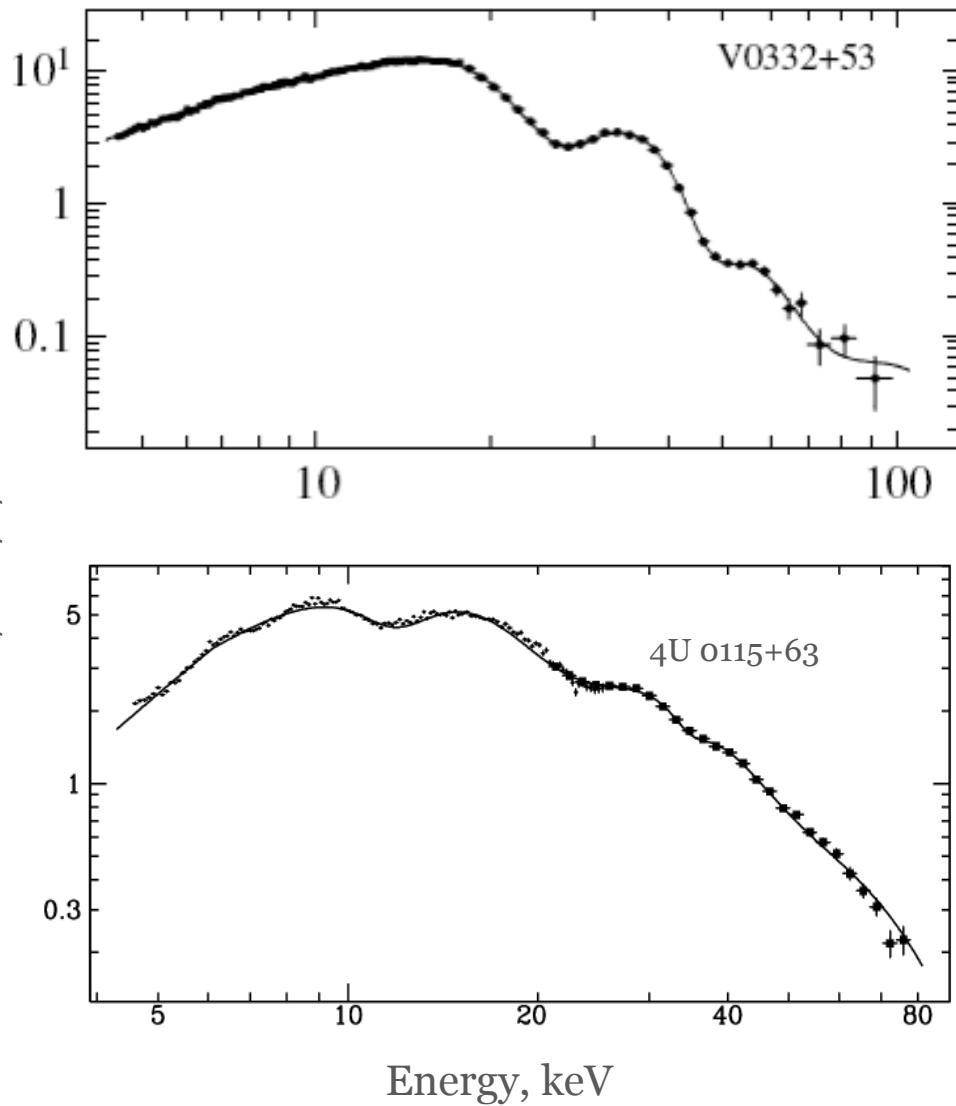
$P_{\text{spin}} \sim 1 - 10^3 \text{ s}$

$B_{NS} \sim 10^{12} \text{ G}$



X-ray pulsar

Typical spectra

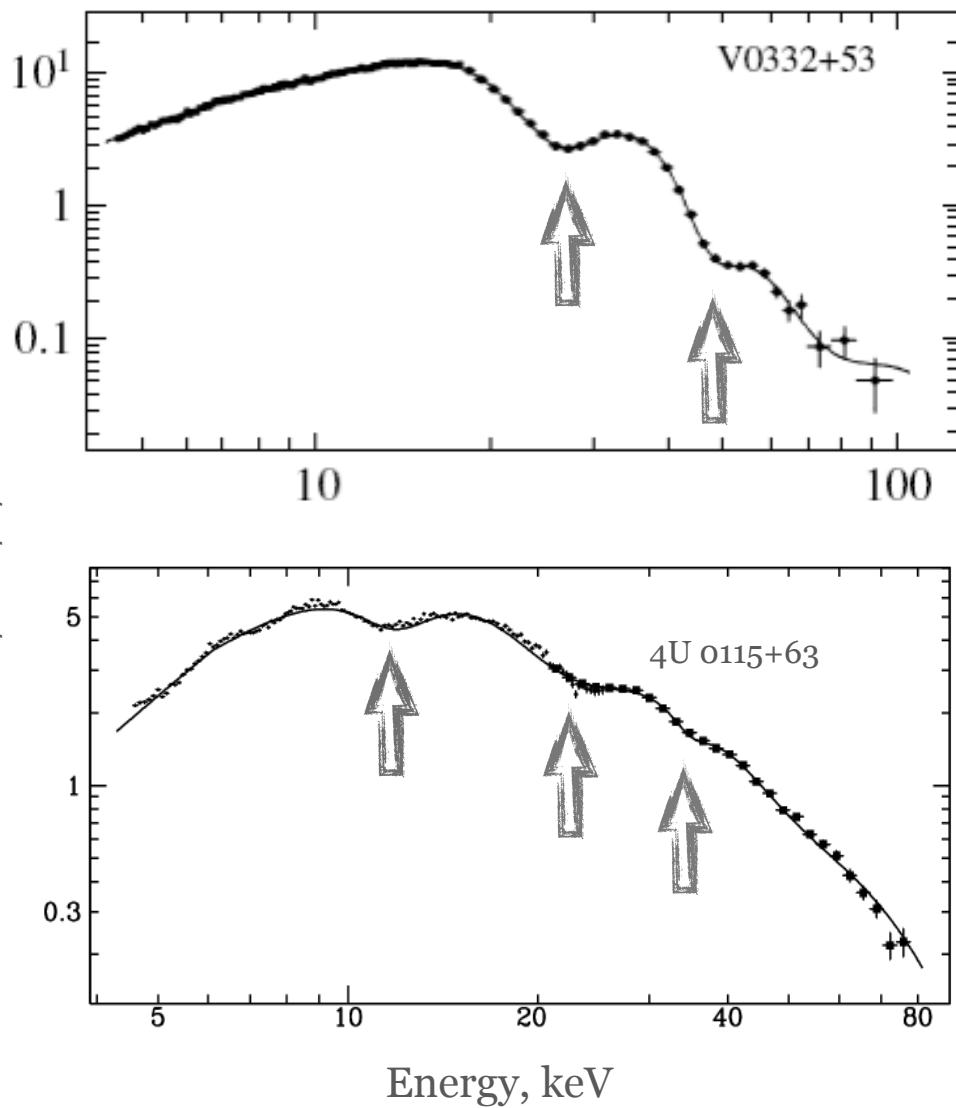


$$E_{\text{cyc}} = 11.6 B_{12} \text{ keV}$$

Source name	Cyclotron energy, keV
4U 0115+63 (-)	11.5, 20.1, 33.6, 49.5, 53
V 0332+53 (-)	28, 53, 74
4U 0352+309 (X Per)	29
RX J0440.9+4431	32
RX J0520.5-6932	31.5
A 0535+262	50, 110
MXB 0656-072	36
Vela X-1 (+)	27, 54
GRO J1008-57	88 [?] , 75.5
1A 1118-61	55
Cen X-3	28
GX 301-2	37, 48
GX 304-1 (+)	50.8
4U 1538-52	20, 47
Swift J1626.6-5156	10
4U 1626-67	37
Her X-1 (+)	42
OAO 1657-415	36
GRO J1744-28	4.7
IGR J18179-1621	21
GS 1843+00	20
4U 1907+09	19, 40
4U 1909+07	44 [?]
XTE J1946+274	36
KS 1947+300	12.5
EXO 2030+375	11 [?] , 36 [?] , 63 [?]
Cep X-4	30

X-ray pulsar

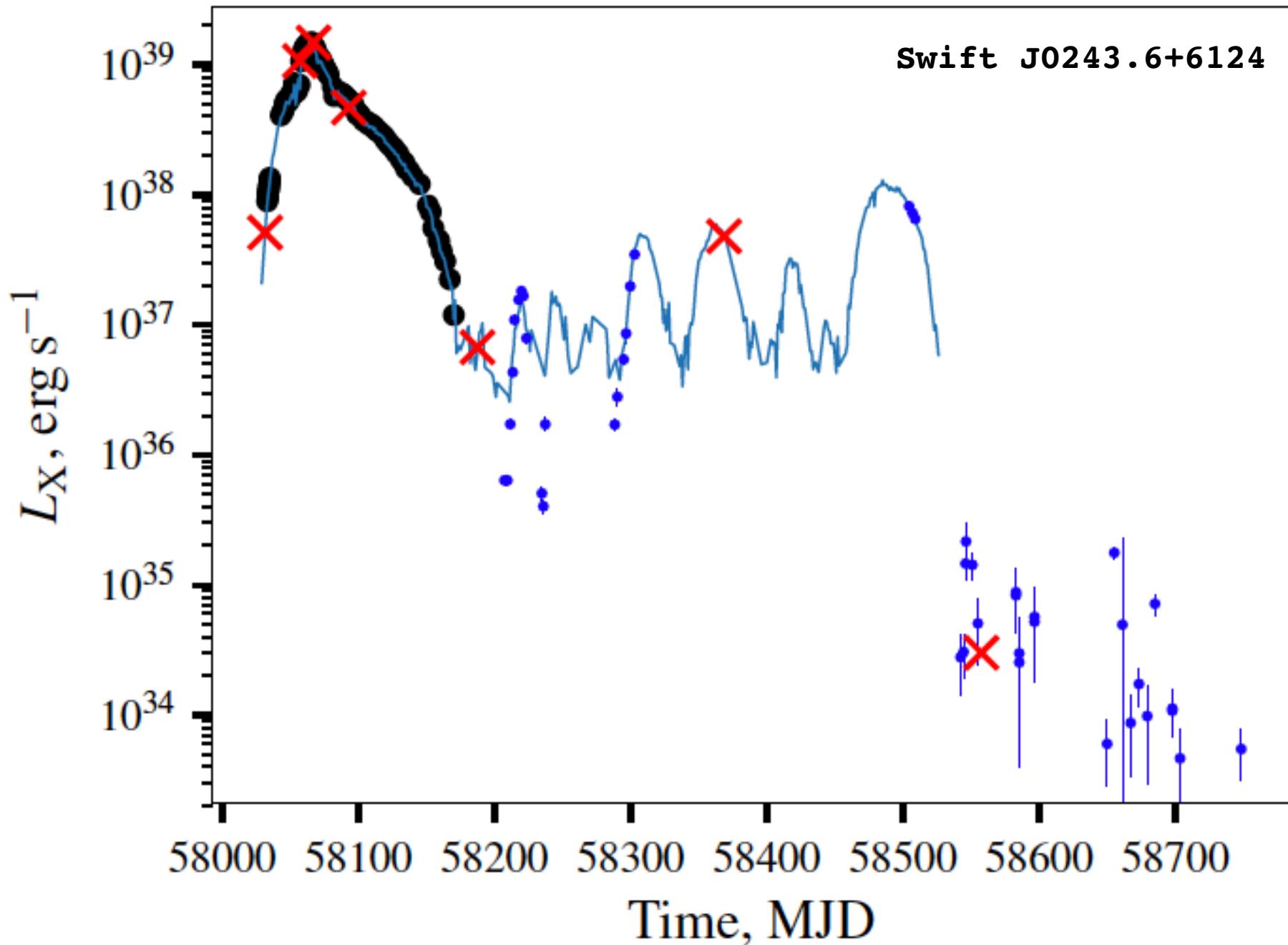
Typical spectra

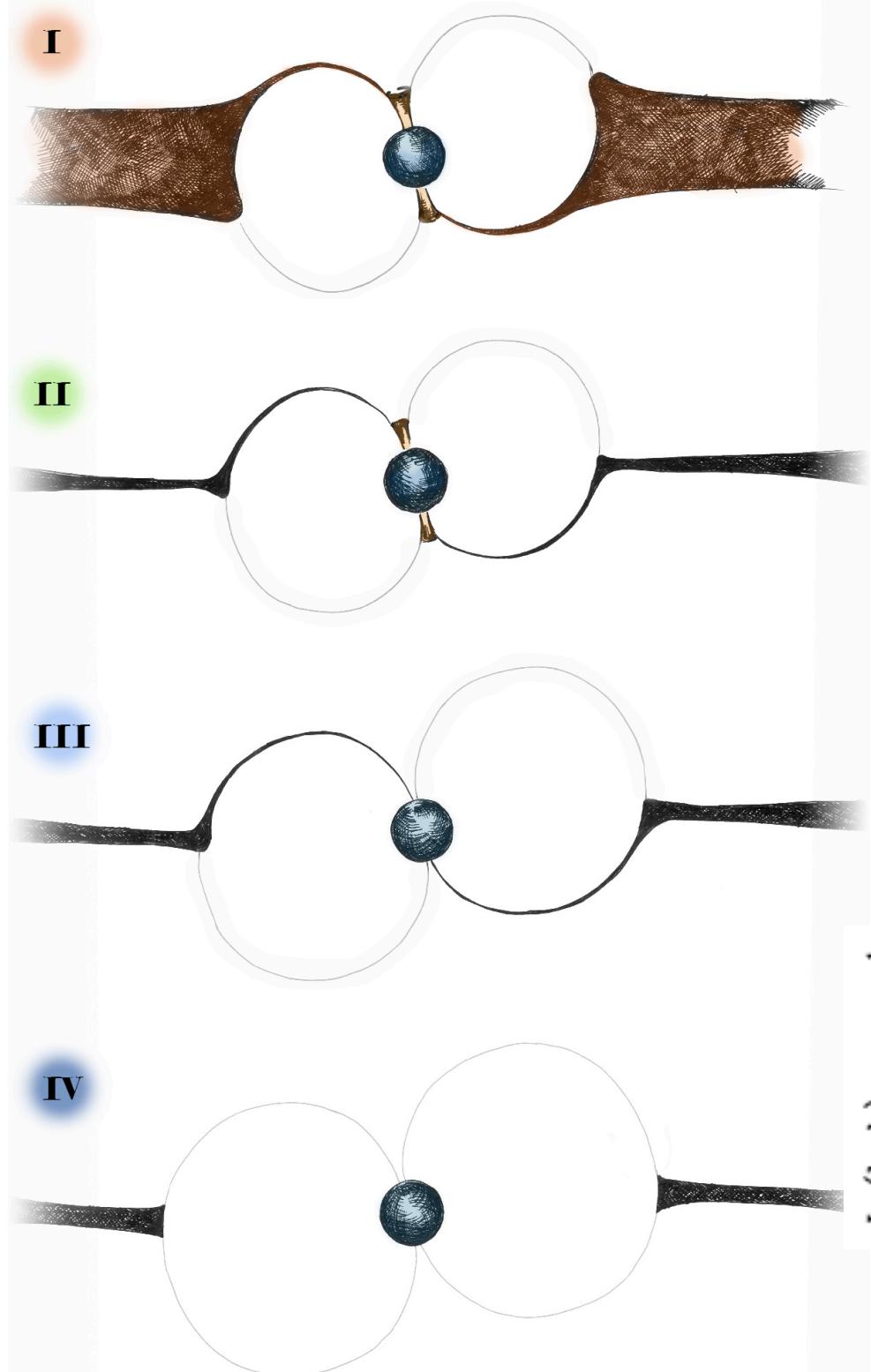


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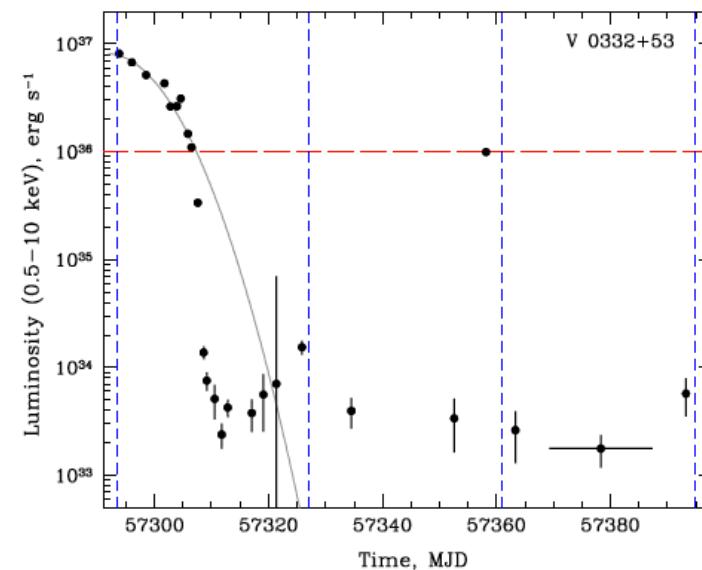
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4U 1909+07	44?
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Cep X-4	30

Transients

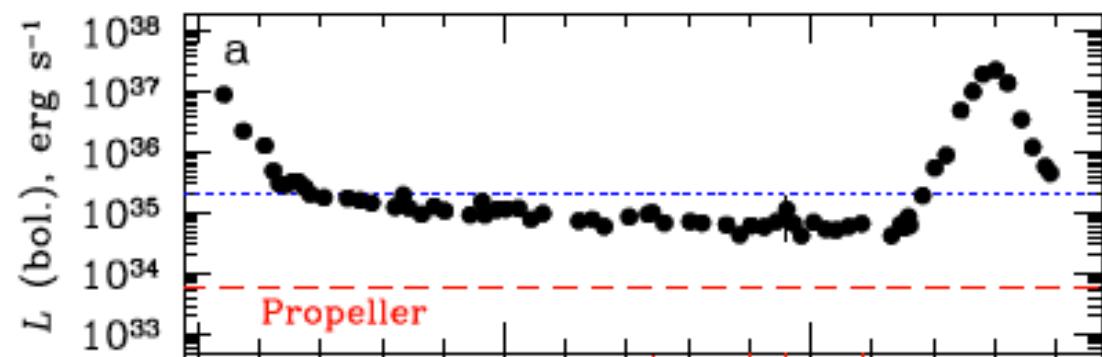




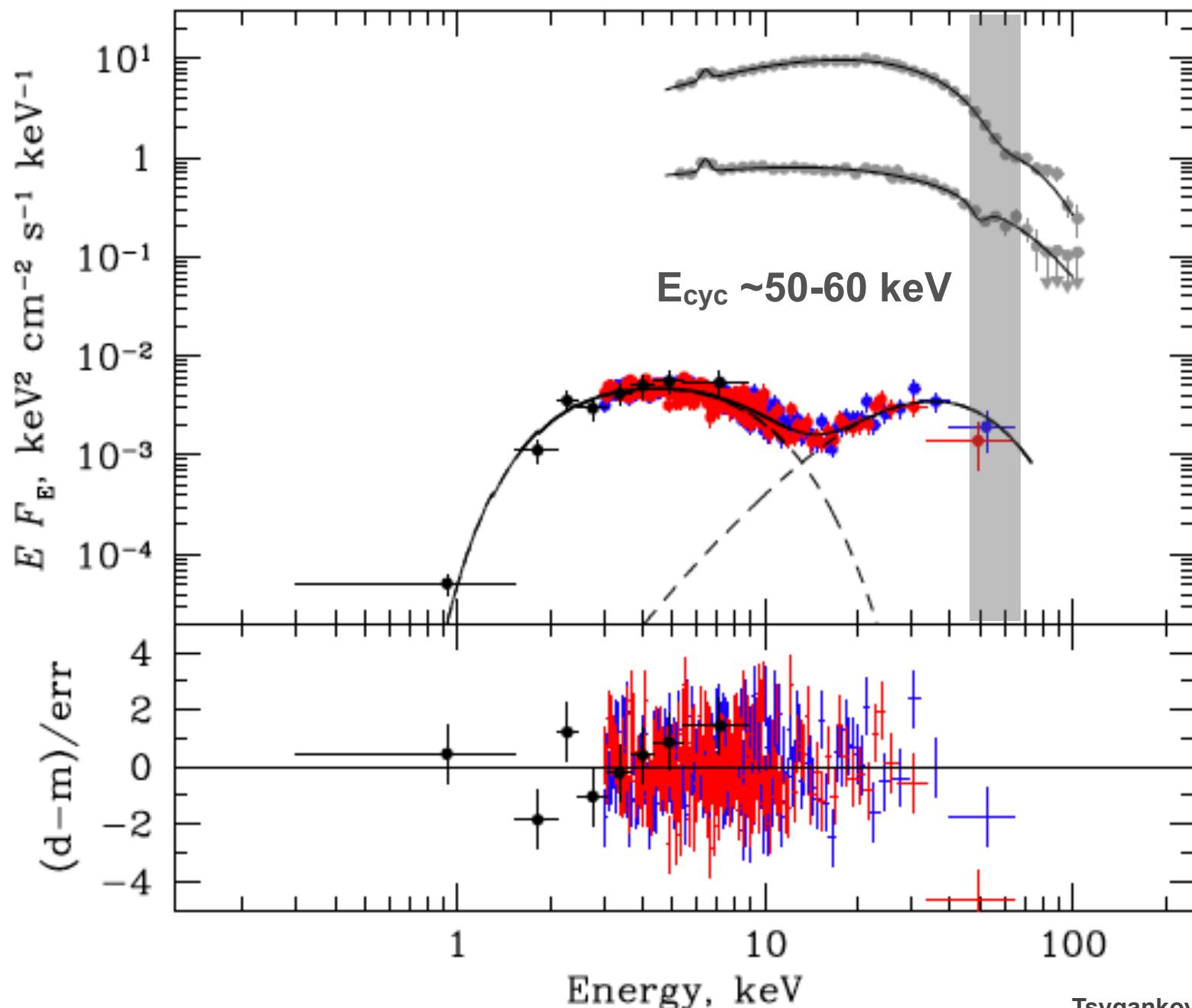
What happens at low luminosity state?



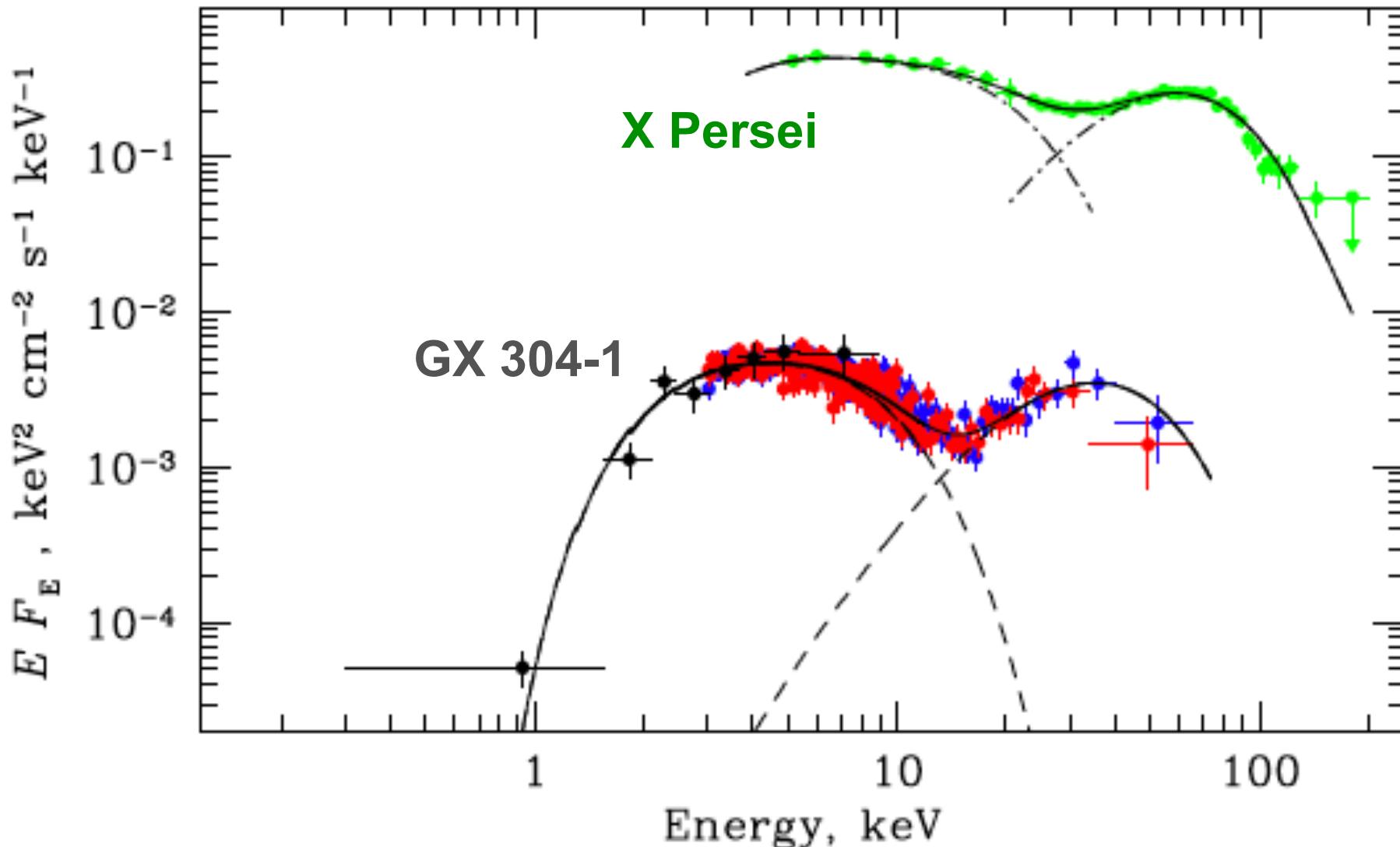
“propeller” effect
vs.
accretion from cold disc



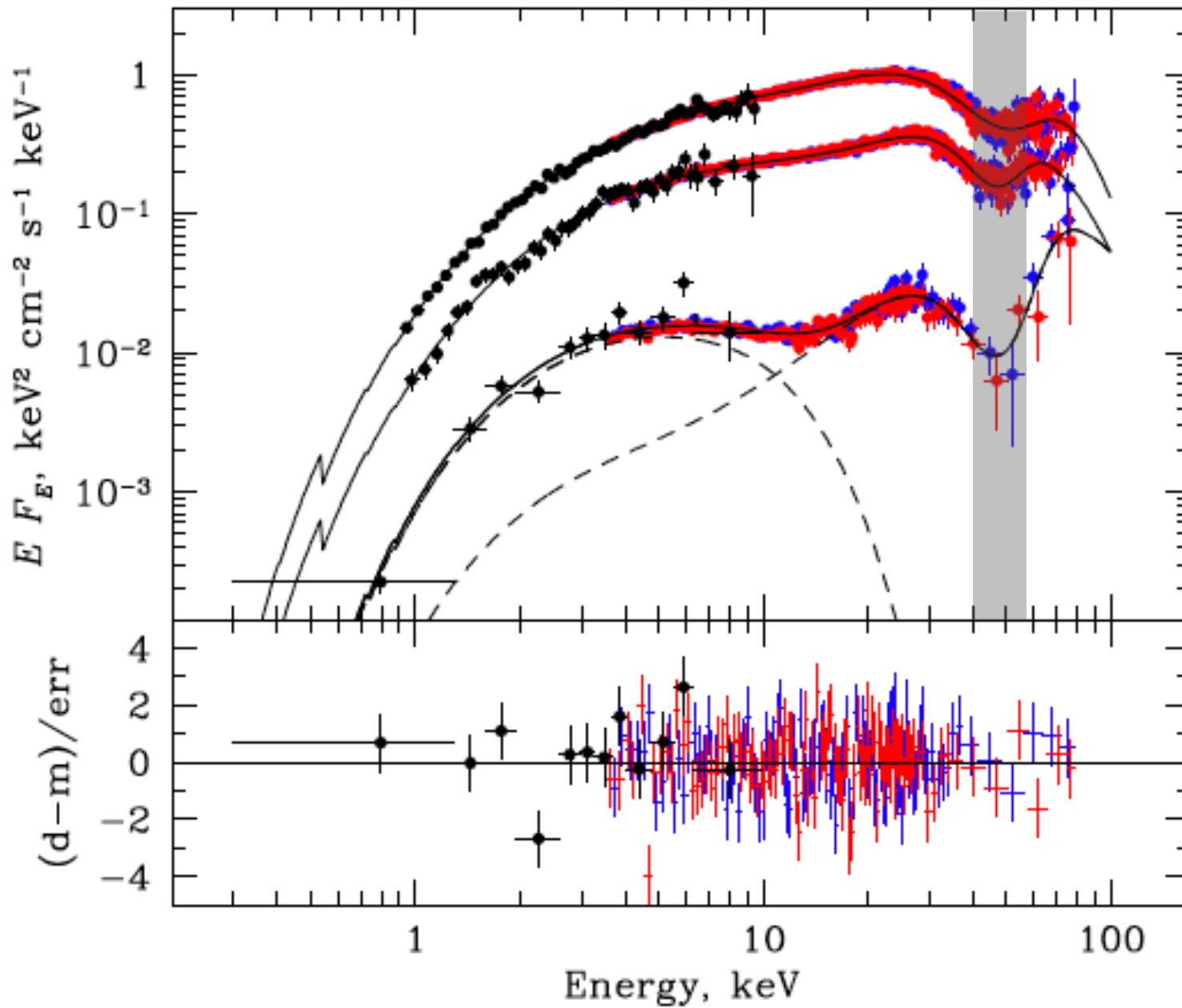
Observations: Low luminosity state in GX 304-1



Observations: Low luminosity state in GX 304-1



Observations: Low luminosity state in A 0535+262

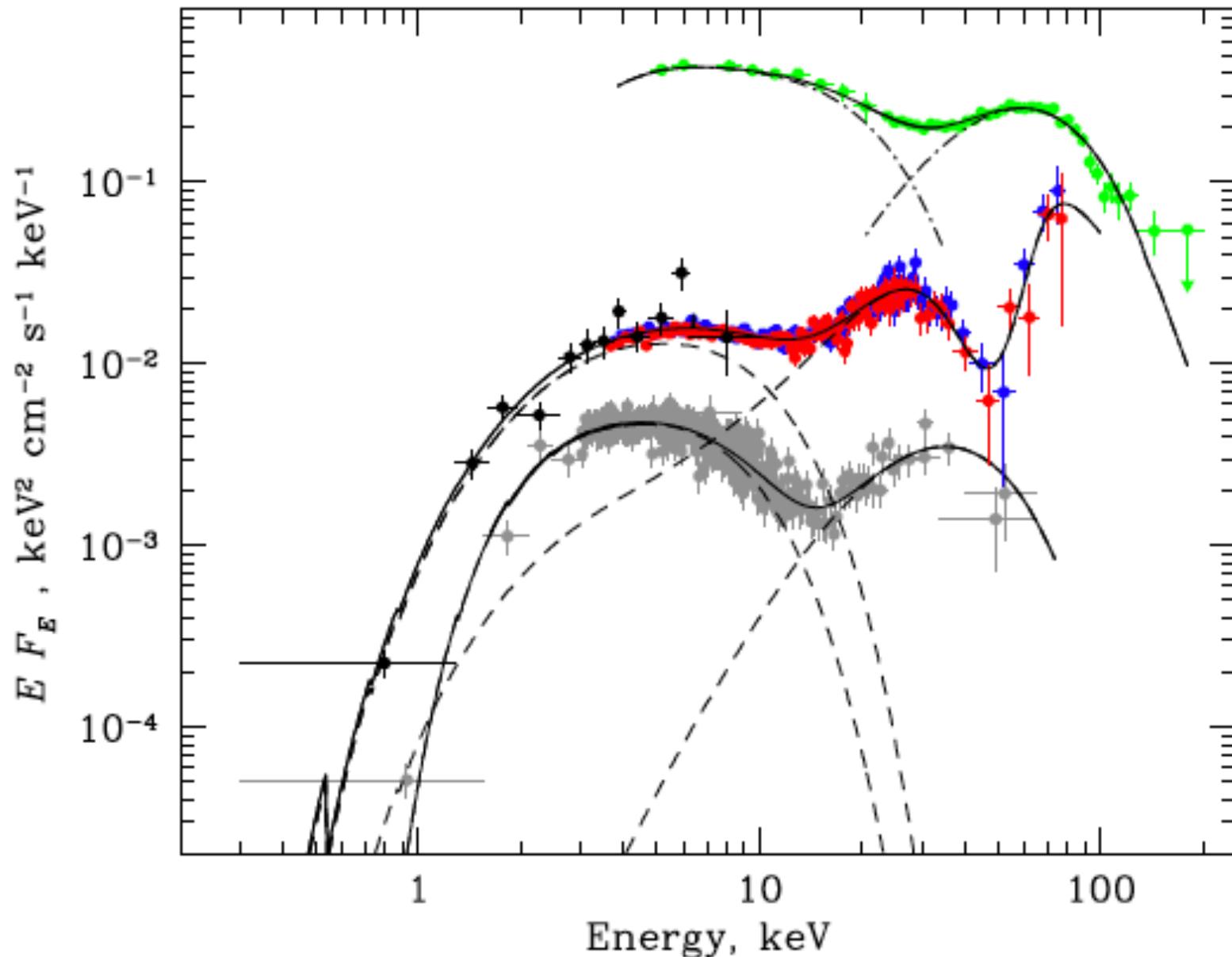


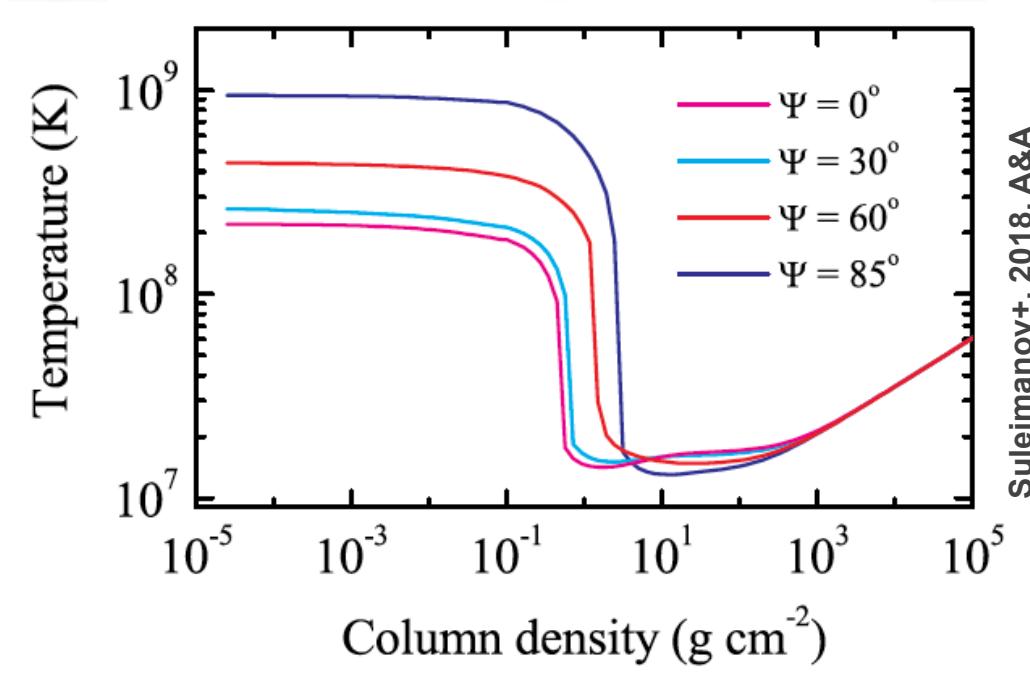
$L \sim 2 \times 10^{36} \text{ erg s}^{-1}$

$L \sim 6 \times 10^{35} \text{ erg s}^{-1}$

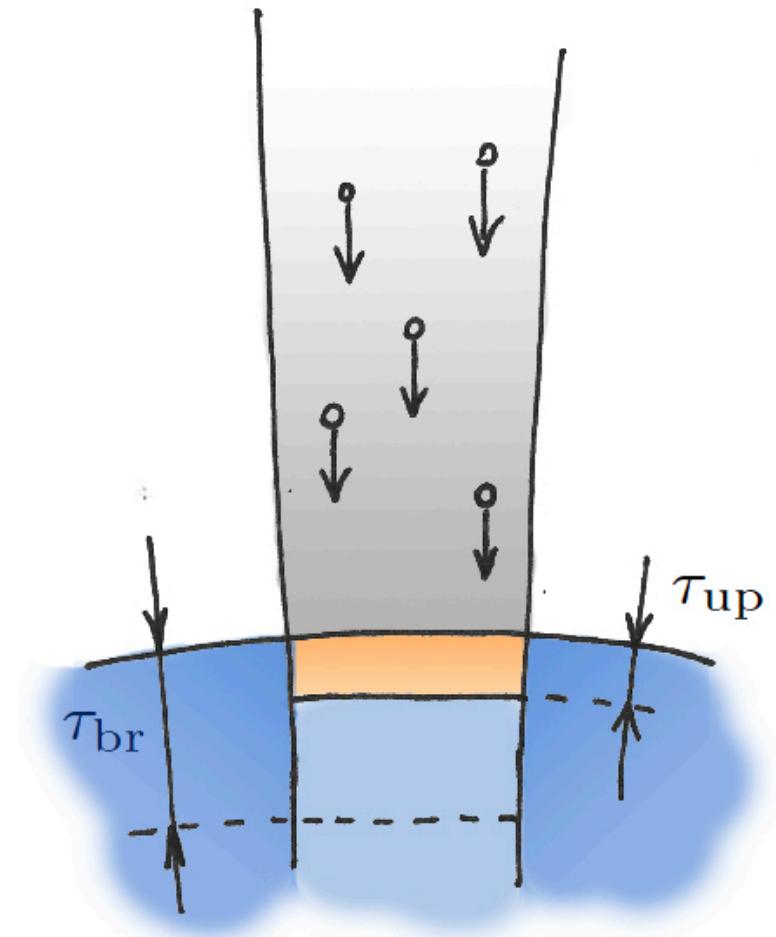
$L \sim 7 \times 10^{34} \text{ erg s}^{-1}$

Observations: Low luminosity state in X Persei, A0535+262 & GX 304-1





What happens at low luminosity state?



Physical model of spectra formation

$$\begin{aligned}\cos \theta \frac{dI_E^{(j)}(\Omega)}{d\tau_T} &= -\frac{\alpha_{\text{abs}}^{(j)}(\Omega)}{\alpha_T} I_E^{(j)}(\Omega) \\ &+ \sum_{i=1}^2 \int_0^\infty dE' \int \frac{d\Omega'}{(4\pi)} \left[R_{ji}(E, \Omega | E', \Omega') I_{E'}^{(i)}(\Omega') - R_{ij}(E', \Omega' | E, \Omega) I_E^{(j)}(\Omega) \right] \\ &+ \frac{\alpha_{\text{abs}}^{(j)}(E, \Omega)}{\alpha_T} \frac{B_E}{2} + S_{\text{ini}}^{(j)}(E, \Omega)\end{aligned}$$

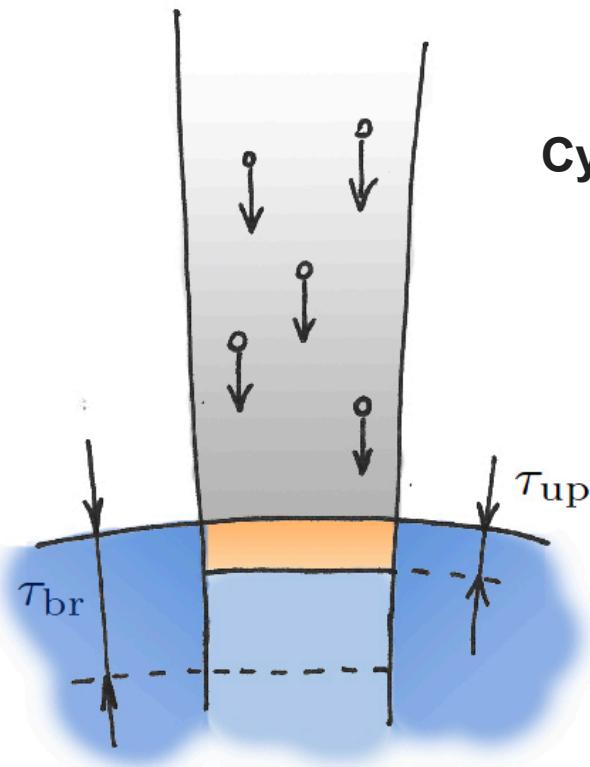
Absorption (bremsstrahlung and cyclotron absorption followed by electron collision)

+

Compton scattering

+

Cyclotron emission and Black Body



Physical model of spectra formation

$$\cos \theta \frac{dI_E^{(j)}(\Omega)}{d\tau_T} = -\frac{\alpha_{\text{abs}}^{(j)}(\Omega)}{\alpha_T} I_E^{(j)}(\Omega) + \sum_{i=1}^2 \int_0^\infty dE' \int \frac{d\Omega'}{(4\pi)} [R_{ji}(E, \Omega | E', \Omega') I_{E'}^{(i)}(\Omega') - R_{ij}(E', \Omega' | E, \Omega) I_E^{(j)}(\Omega)] + \frac{\alpha_{\text{abs}}^{(j)}(E, \Omega)}{\alpha_T} \frac{B_E}{2} + S_{\text{ini}}^{(j)}(E, \Omega)$$

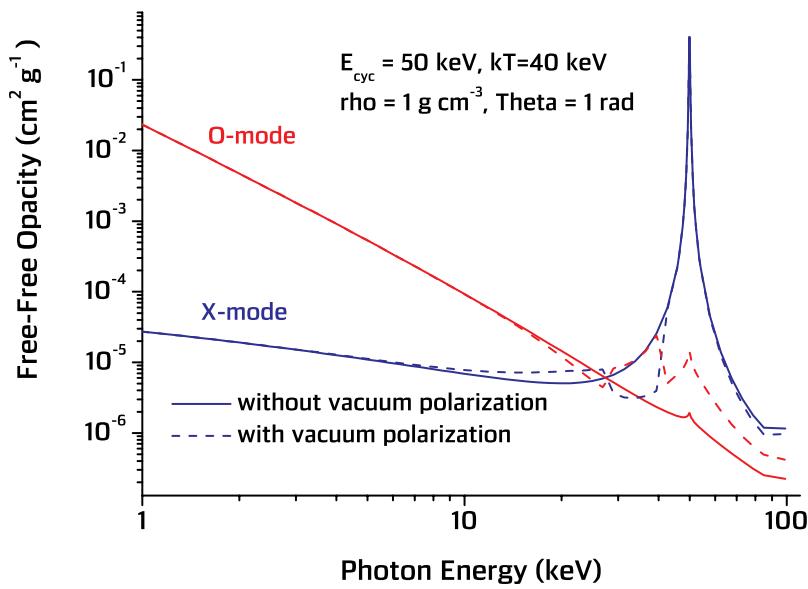
Absorption (bremsstrahlung and cyclotron absorption followed by electron collision)

+

Compton scattering

+

Cyclotron emission and Black Body



Probability of true cyclotron absorption:

$$P_{\text{abs, true}} \simeq \frac{r_{\text{coll}}}{r_{\text{cyc}}} \sim 1.7 \times 10^{-7} n_{e,21} B_{12}^{-7/2}$$

Physical model of spectra formation

$$\begin{aligned}
 \cos \theta \frac{dI_E^{(j)}(\Omega)}{d\tau_T} = & -\frac{\alpha_{\text{abs}}^{(j)}(\Omega)}{\alpha_T} I_E^{(j)}(\Omega) \\
 & + \sum_{i=1}^2 \int_0^\infty dE' \int \frac{d\Omega'}{(4\pi)} \left[R_{ji}(E, \Omega | E', \Omega') I_{E'}^{(i)}(\Omega') - R_{ij}(E', \Omega' | E, \Omega) I_E^{(j)}(\Omega) \right] \\
 & + \frac{\alpha_{\text{abs}}^{(j)}(E, \Omega)}{\alpha_T} \frac{B_E}{2} + S_{\text{ini}}^{(j)}(E, \Omega)
 \end{aligned}$$

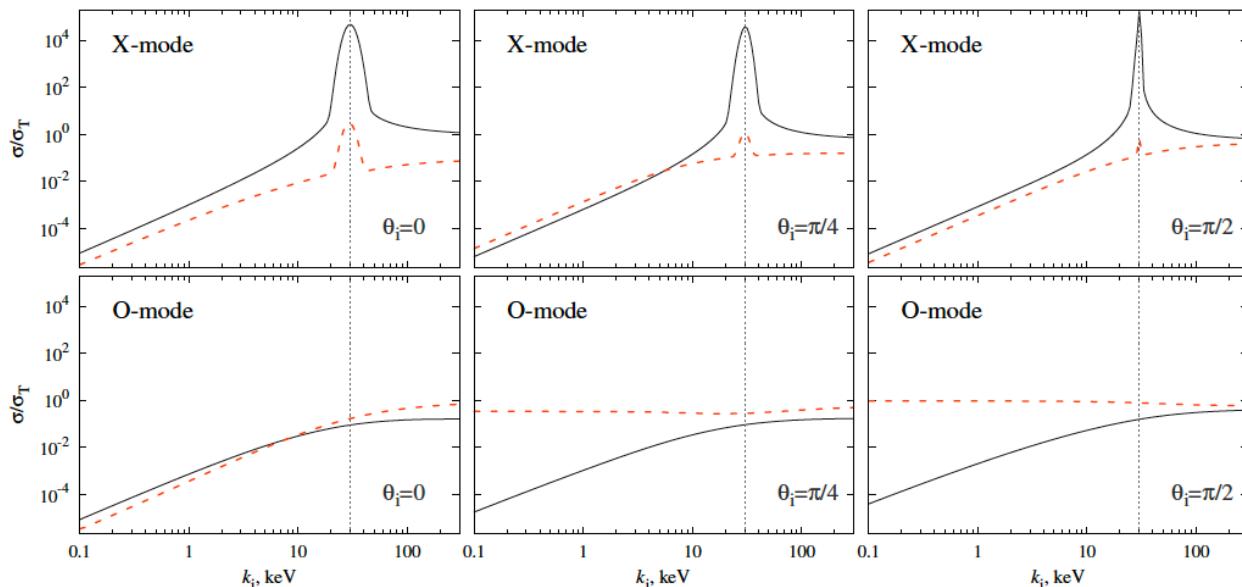
Absorption (bremsstrahlung and cyclotron absorption followed by electron collision)

+

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Cross-sections are based
on non-relativistic amplitudes
recalculated for plasma
polarisation modes

Physical model of spectra formation

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Absorption (bremsstrahlung and cyclotron absorption followed by electron collision)

+

Compton scattering

+

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Details of numerical simulations

The basic parameters in the model:

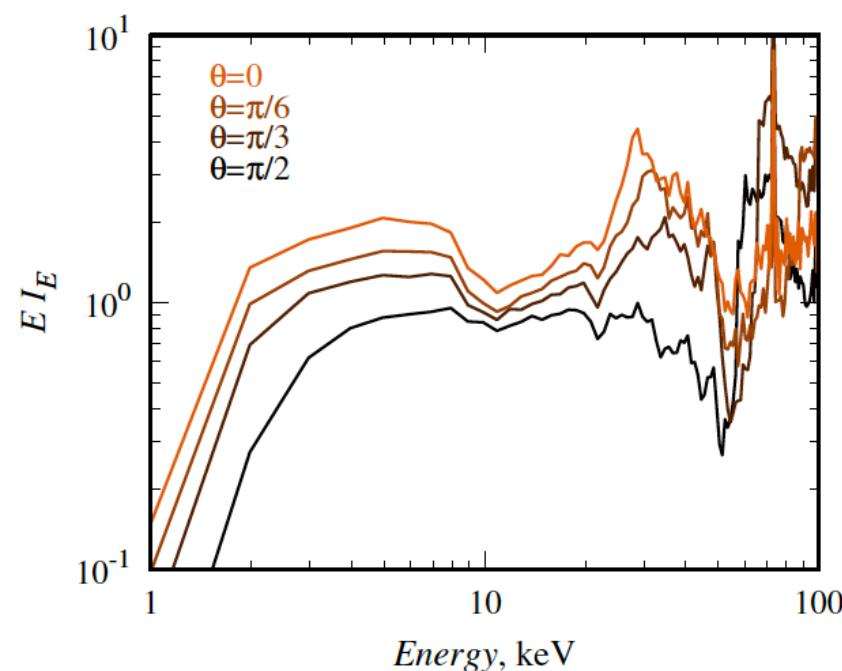
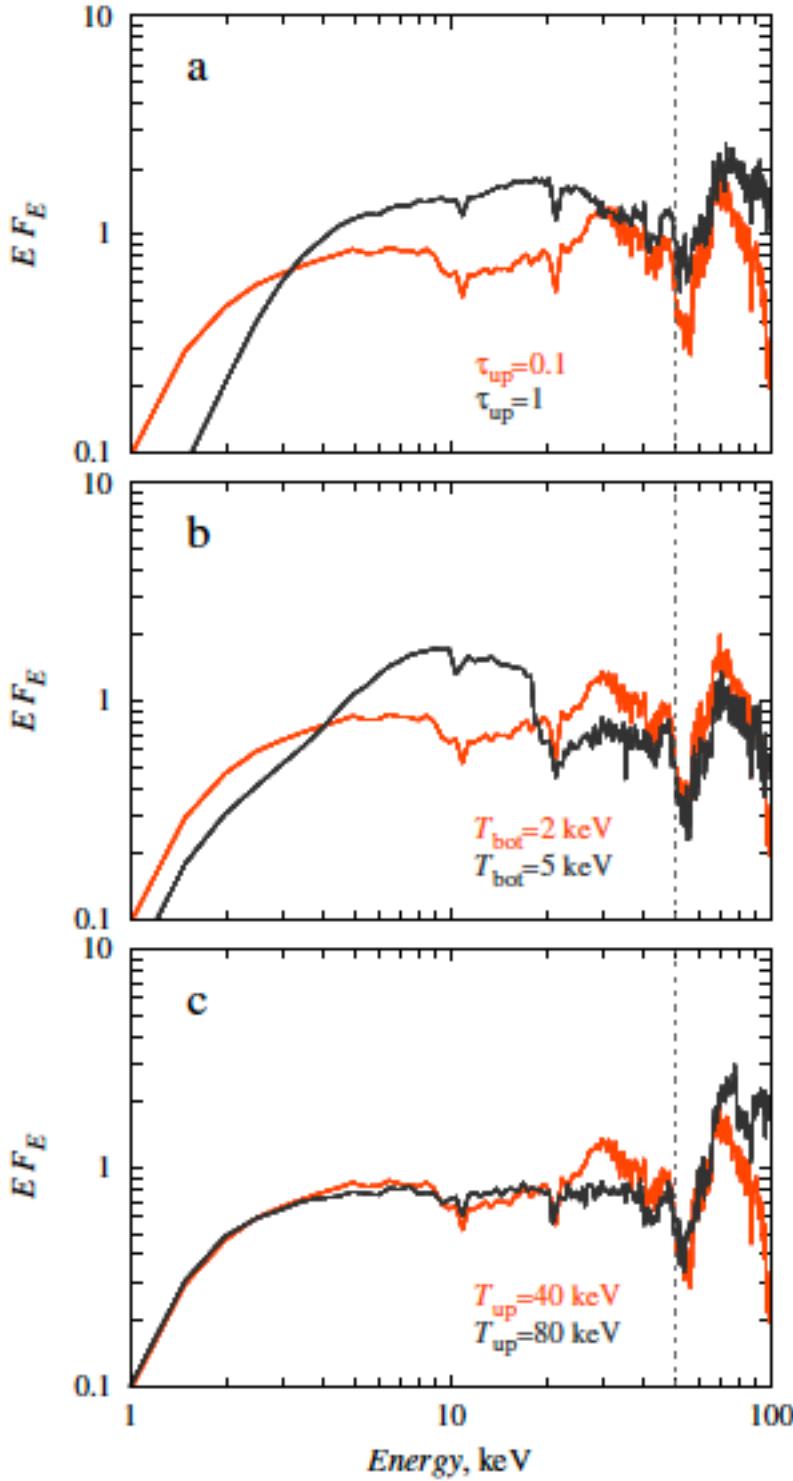
- (1) temperature of the upper layer (T_{up}) and temperature below the upper layer (T_{bot});
- (2) optical thickness of the upper layer due to Thomson scattering (τ_{up});
- (3) effective depth where the accretion flow stopes in the atmosphere (τ_{br}).

$$E_{\text{cyc}} = 50 \text{ keV}$$

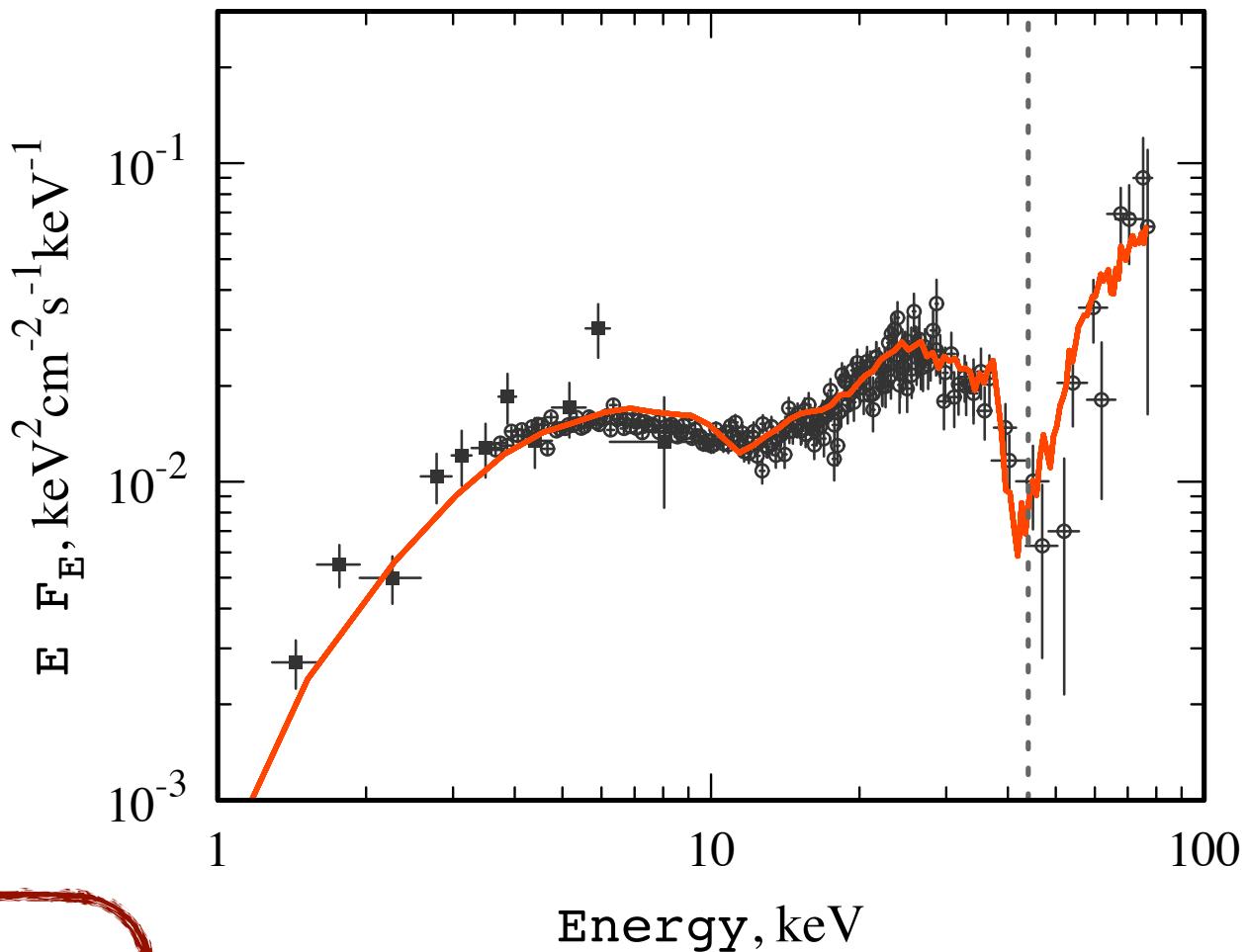
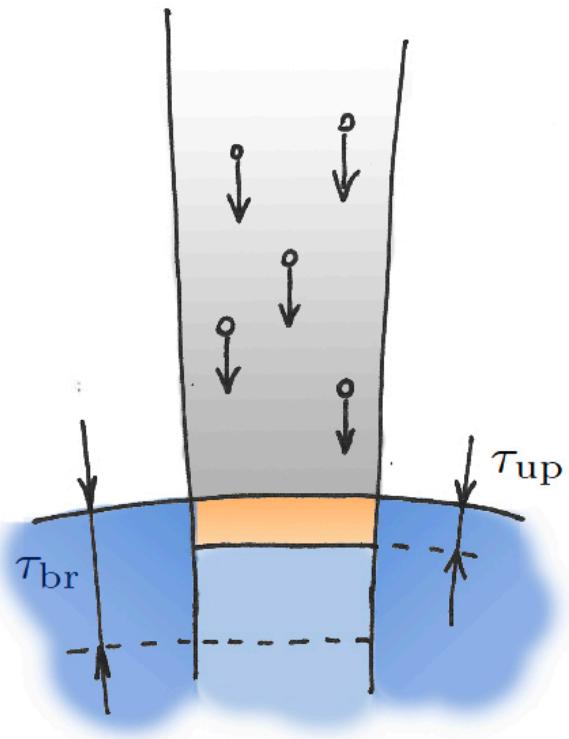
$$T_{\text{bot}} = 2 \text{ keV}, T_{\text{up}} = 40 \text{ keV}$$

$$\tau_{\text{up}} = 0.1, \tau_{\text{br}} = 1$$

$$Z = 1$$



Results of numerical simulations vs. Data



Set of Parameters:

$$\tau_{\text{br}} = 10$$

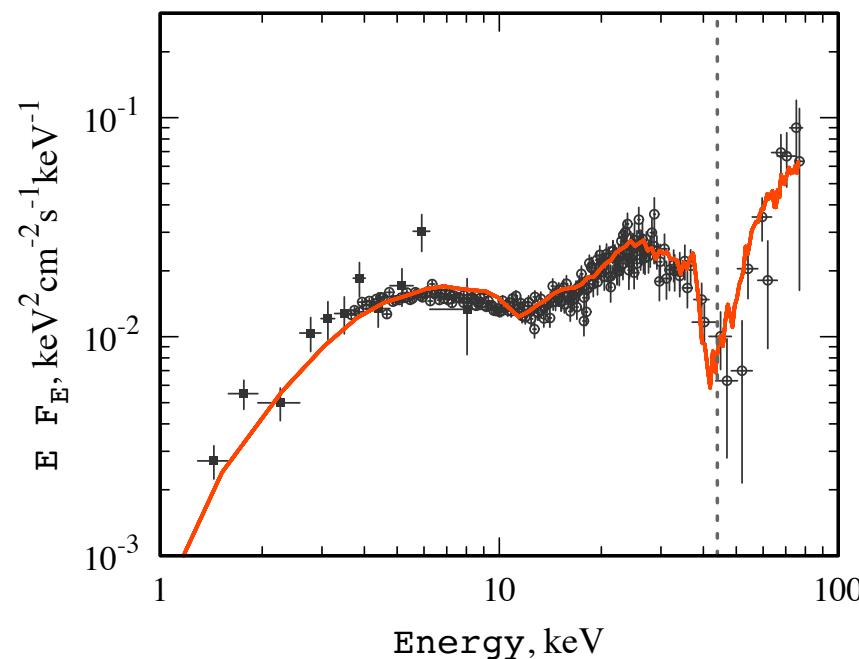
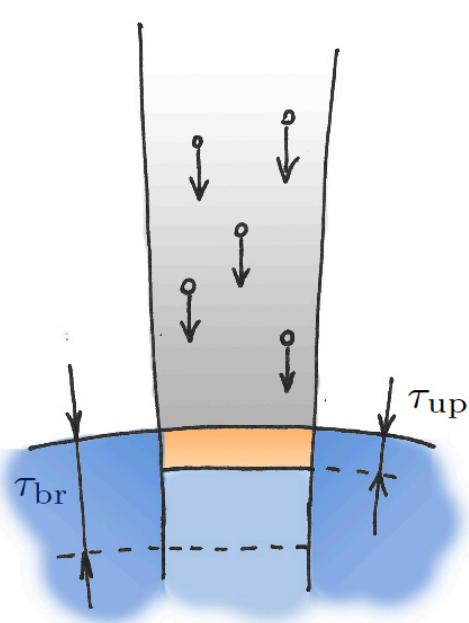
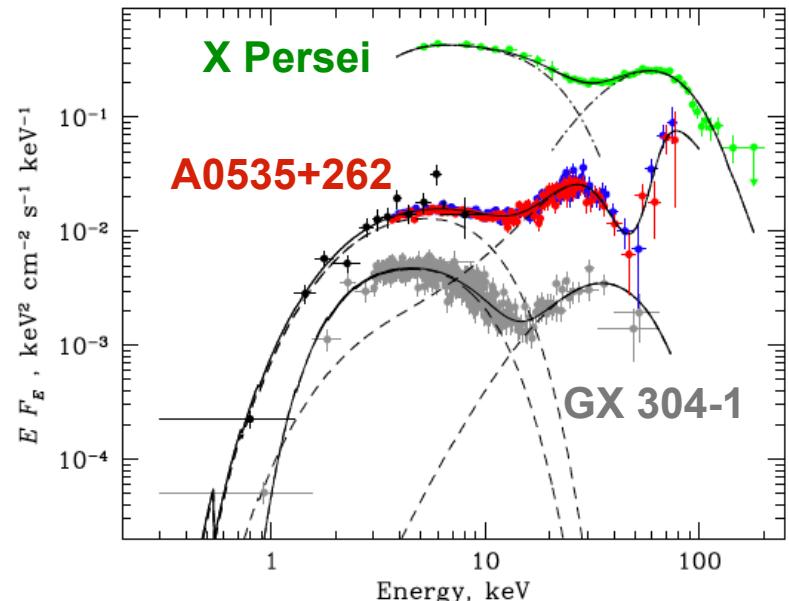
$$\tau_{\text{up}} = 0.5 \quad T_{\text{up}} \approx 100 \text{ keV}$$

$$T_{\text{down}} \approx 2.7 \text{ keV}$$

Conclusions

There is **dramatic changes** of X-ray energy spectra at very low luminosity states of X-ray pulsars. The spectra clearly **show two components**.

The observed spectra can be explained by cyclotron emission in the atmosphere of accreting neutron star with its further comptonisation affected by resonant scattering.



Spectra of X-ray pulsars at very low luminosity states can be used for estimations of magnetic field strength at the neutron star surface.

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